

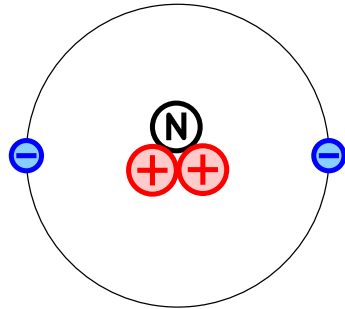
$^3\text{He}$

Superfluidity, collective modes, spin waves

Vladislav Zavjalov

# Helium: two stable isotopes

$^3\text{He}$

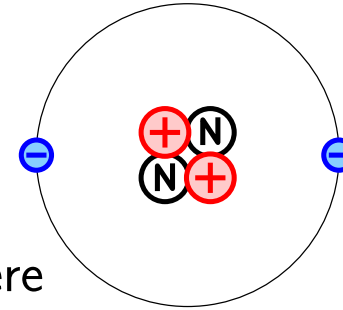


2 protons  
1 neutron  
2 electrons

1:1'000'000 ratio in atmosphere

Spin:  $1/2$ , Fermi particle

$^4\text{He}$



2 protons  
2 neutrons  
2 electrons

Spin: 0, Bose particle

Inert atoms with symmetric electron shells.

Weak interactions between atoms, small atom mass.

Liquid down to absolute zero temperature (at small pressures).

$^3\text{He}$ :

Superfluid transition at  $\sim 1$  mK

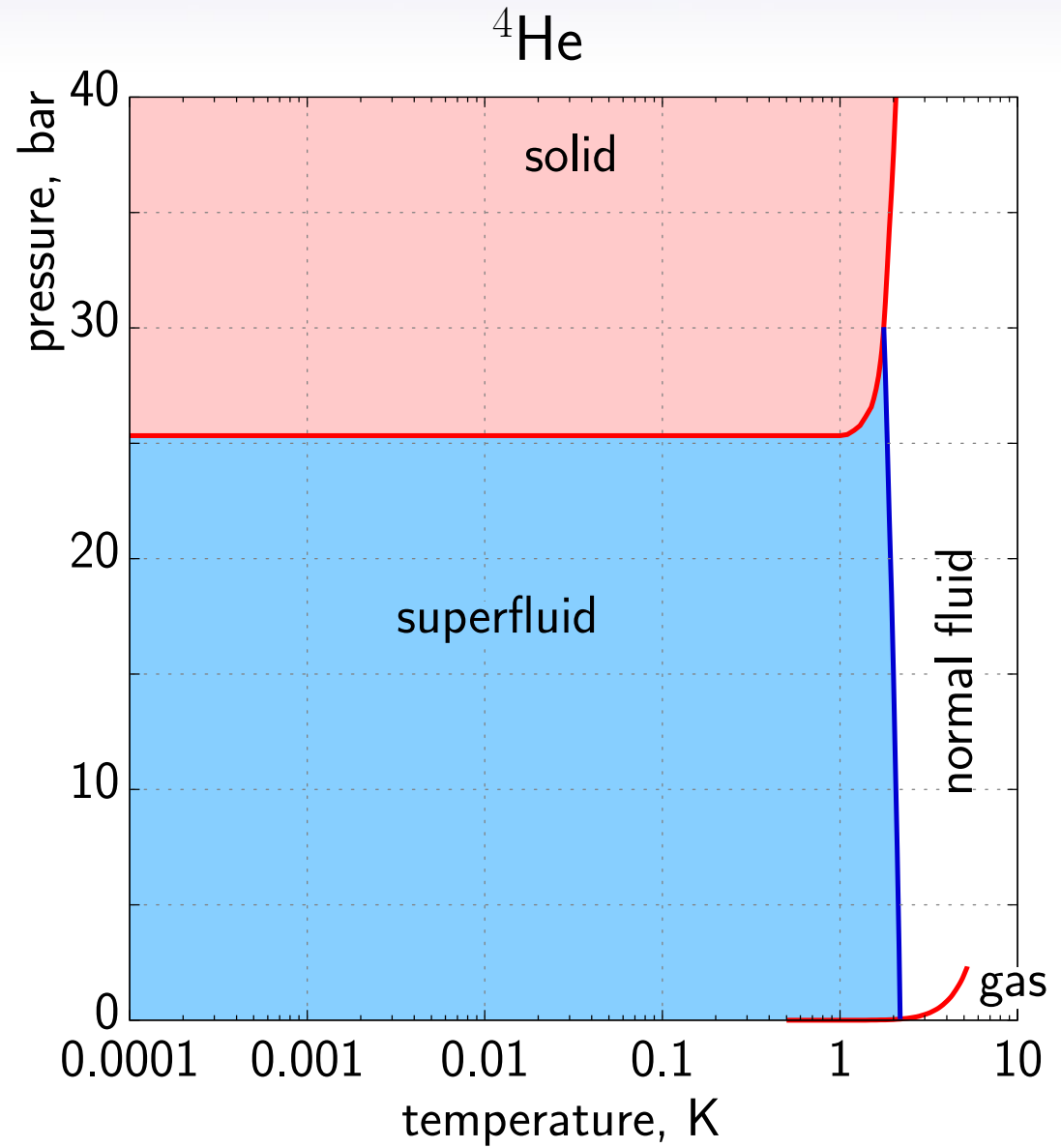
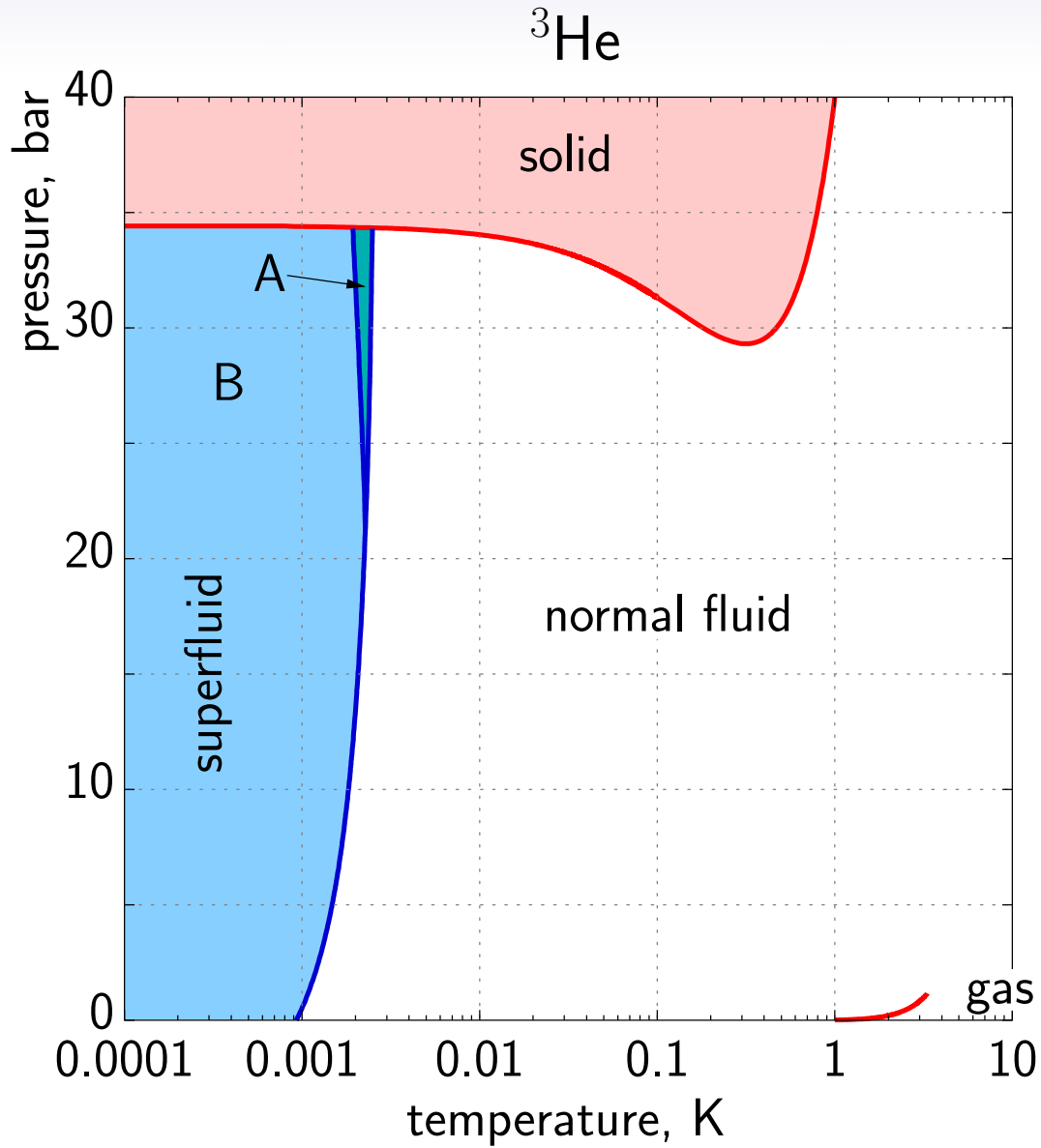
Cooper pairing with  $L = 1$  and  $S = 1$ .

$^4\text{He}$ :

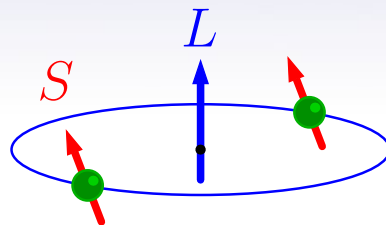
Superfluid transition at  $\sim 2$  K

Bose condensation of atoms.

# Helium: Phase diagram

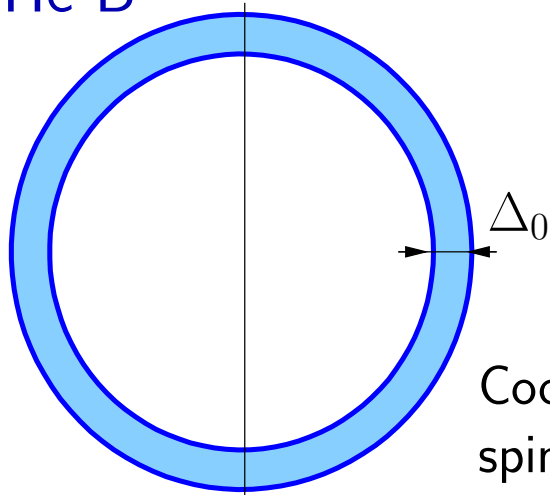


# Superfluid $^3\text{He}$



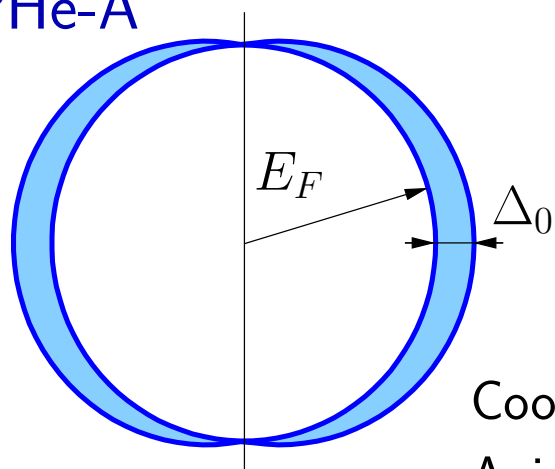
Cooper pair  
 $S = 1, L = 1$

$^3\text{He-B}$

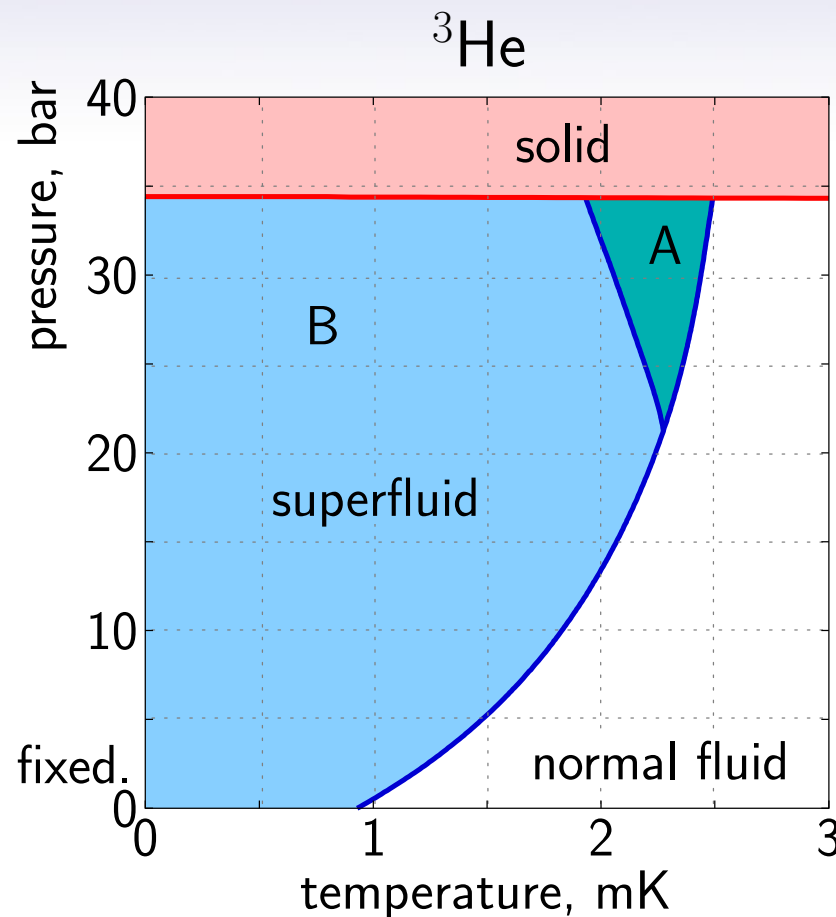


Cooper pairs with all possible spin projections  $S_z = 0, \pm 1$ .  
 Mutual orientation of  $S$  and  $L$  is fixed.  
 Isotropic energy gap.

$^3\text{He-A}$



Cooper pairs with  $S_z = \pm 1$ .  
 Anisotropic energy gap.



Superfluid  $^3\text{He}$ :

- complicated theory
- pure system, good agreement with experiments
- macroscopic system
- no applications

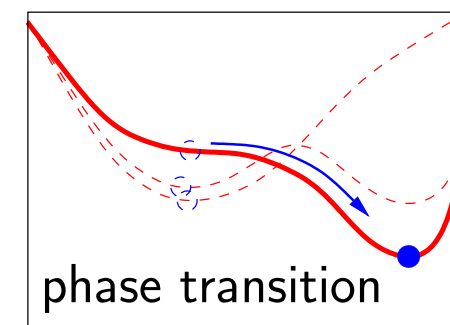
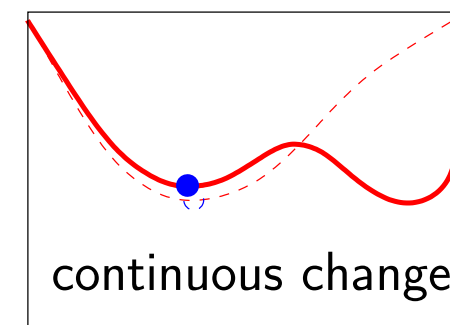
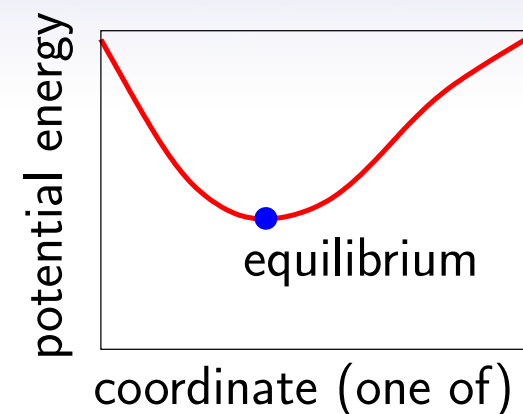
Phase transition: sudden jump of a system to a new state during continuous change of some external condition (temperature, pressure, magnetic field...).

A system state can be represented by a point in a multi-dimensional coordinate space.

Potential energy depends on coordinates, it continuously changes following external conditions.

First order phase transition:

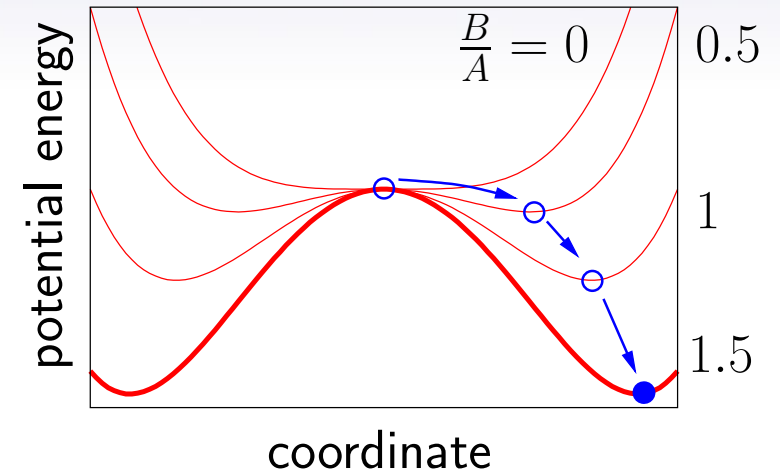
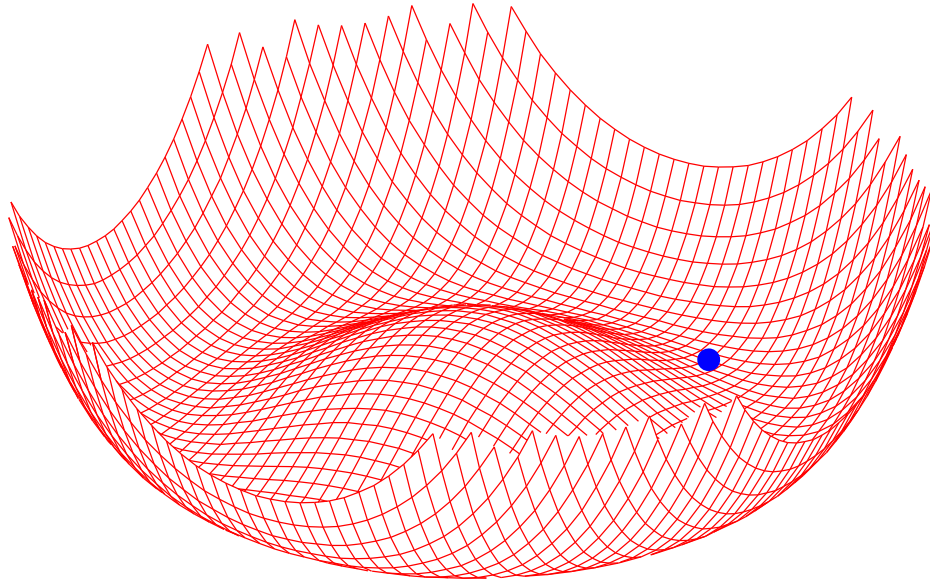
- All properties of the system have a jump.
- Hysteresis: backward transition at a different point.
- Latent heat of the transition, singularity in heat capacity (extra kinetic energy has to be removed or added to thermalize the system).



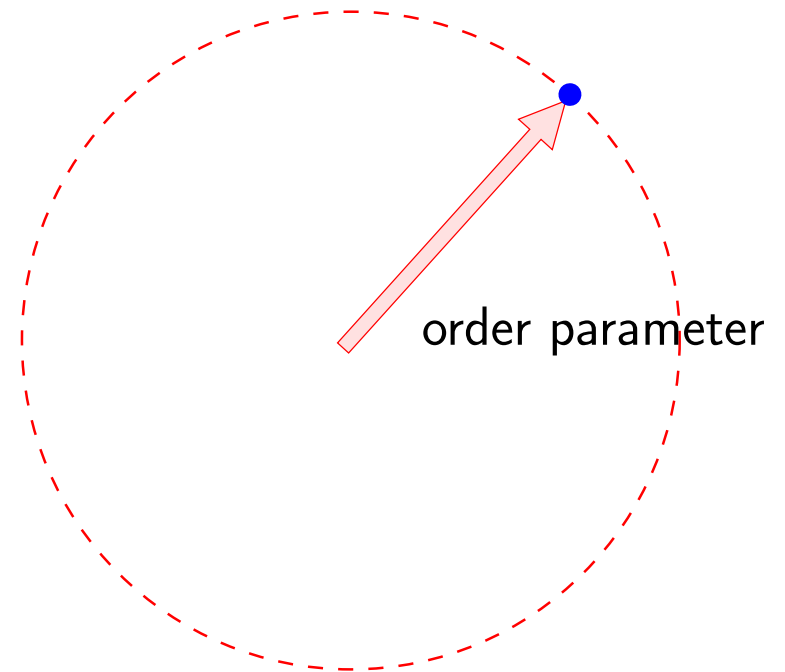
# Phase transitions: 2nd order

“Mexican-hat”, or “wine-bottle” potential.

Ginzburg-Landau model:  $A r^4 - B r^2$



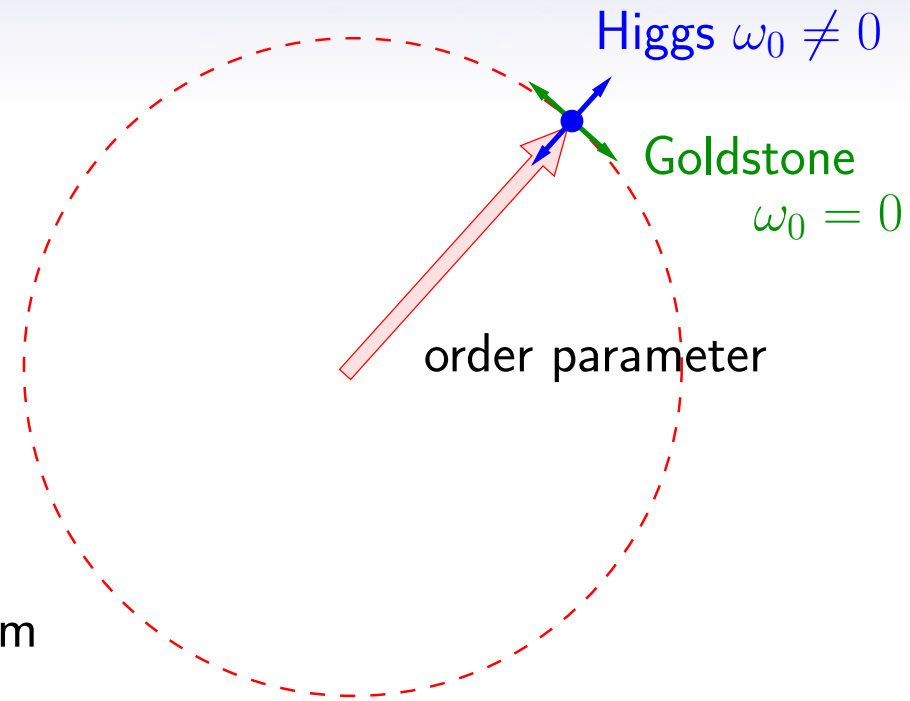
- System state changes continuously, but a new degree of freedom appears.
- No hysteresis.
- Jump in heat capacity.
- Spontaneous symmetry breaking
- Order parameter (phase and amplitude).



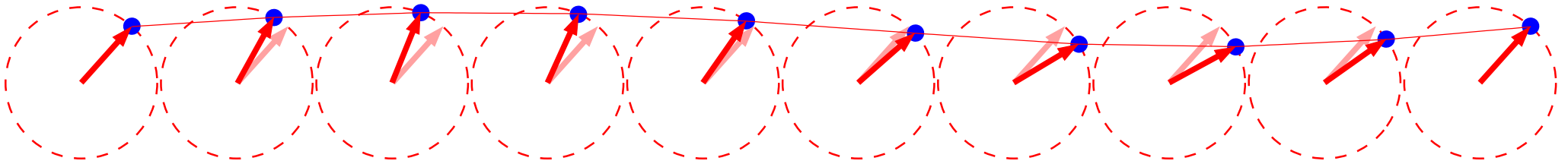
Order parameter space: 2D plane  
2 degrees of freedom,  
2 oscillation modes.

Degenerate space: a circle where energy is at minimum  
1 degree of freedom,  
1 Goldstone mode - motion of the order parameter phase

Higgs mode - motion of the order parameter amplitude.  
non-zero frequency  $\omega_0$ .



We can have a Goldstone wave with non-zero frequency:



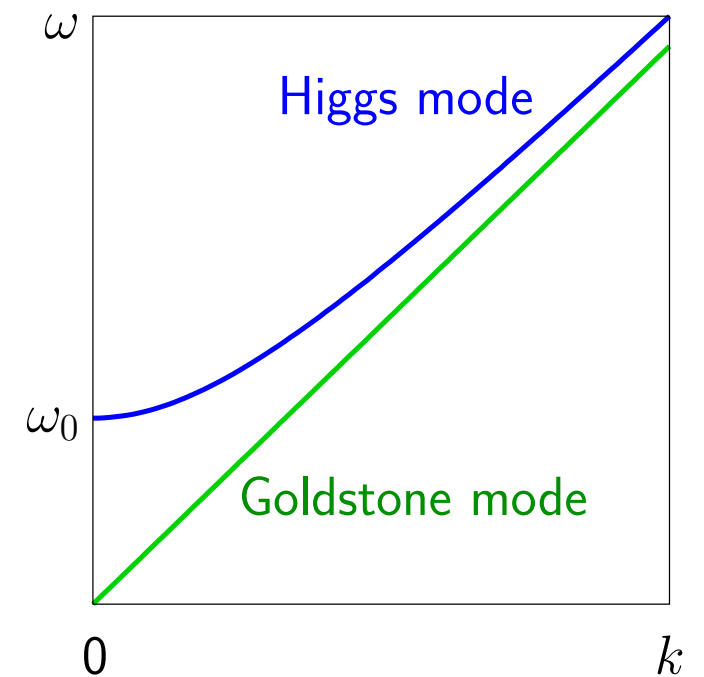
Wave length  $L$ , wave vector  $k = 2\pi/L$

$\omega = kc$  where  $c$  is speed of the wave

For Higgs mode minimum frequency is  $\omega_0$

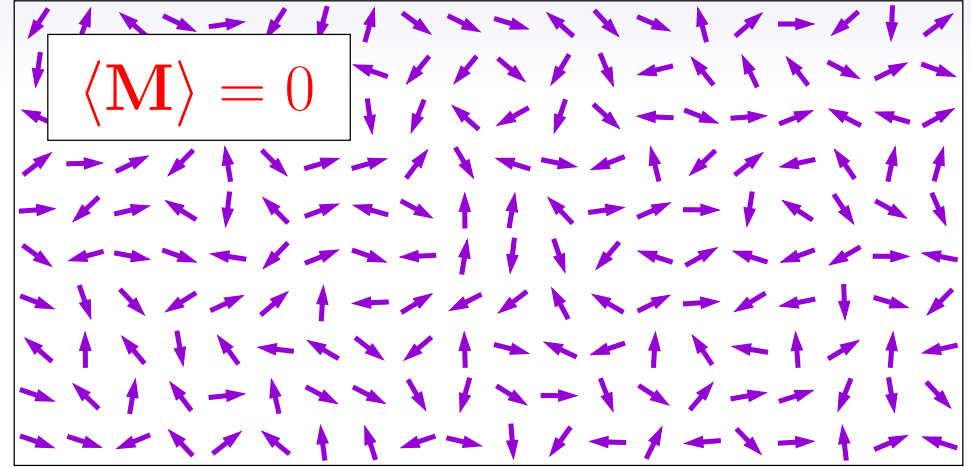
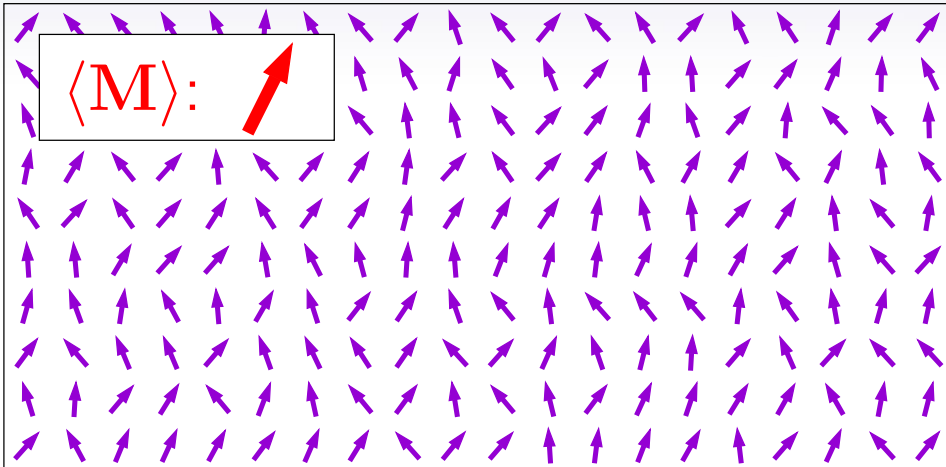
$$\omega = \sqrt{\omega_0^2 + (ck)^2}$$

Relativistic spectrum with non-zero rest mass.





# Simple example: ferromagnet



—————→ temperature

## Order parameter - magnetization $\langle \mathbf{M} \rangle$

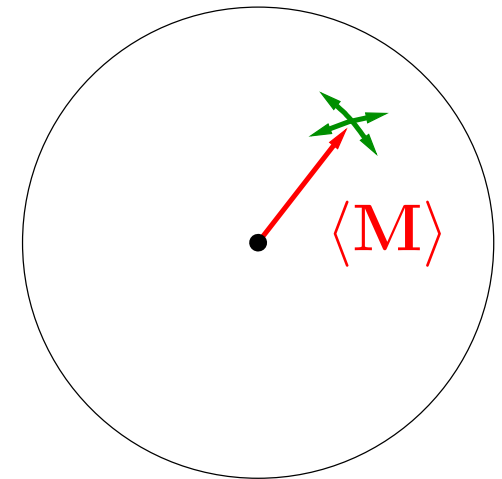
Order parameter space: 3D  $\rightarrow$  three modes.

Degenerate space: a sphere  $\rightarrow$  two Goldstone modes.

Higgs mode: oscillation of magnetization amplitude

Goldstone modes: rotation of magnetization, spin waves

(2 modes, can be splitted by magnetic field)



<sup>4</sup>He (and superconductors):

Order parameter:  $A = \Delta e^{i\varphi}$  (wave function of Bose condensate, a complex number)

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<sup>3</sup>He:

Cooper pairs have  $S = 1$  and  $L = 1$ , the condensate state is a linear combination of states with  $L_z = -1, 0, +1$  and  $S_z = -1, 0, 1$

Order parameter:  $A_{\mu j}$ , a 3x3 complex matrix. 18 degrees of freedom!

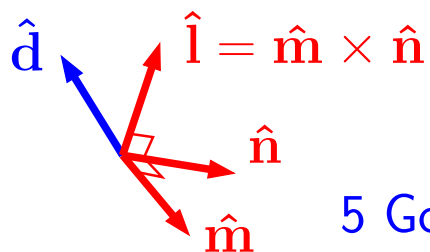
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Symmetry breaks in different ways in A- and B-phases:

<sup>3</sup>He-A

$$A_{\mu j} = \Delta_0 \hat{\mathbf{d}}_{\mu} (\hat{\mathbf{m}}_j + i \hat{\mathbf{n}}_j)$$

$\hat{\mathbf{m}}$  and  $\hat{\mathbf{n}}$  - orthogonal unit vectors

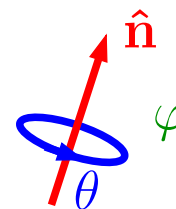


5 Goldstone modes

<sup>3</sup>He-B

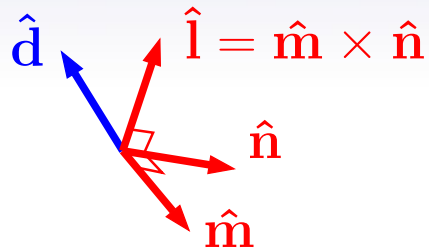
$$A_{\mu j} = \Delta_0 R_{\mu j}(\hat{\mathbf{n}}, \theta) e^{i\varphi}$$

$R_{\mu j}$  - rotation matrix with axis  $\hat{\mathbf{n}}$  and angle  $\theta$



4 Goldstone modes

$^3\text{He-A}$



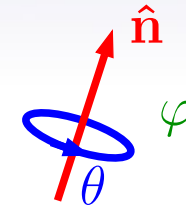
Goldstone modes (5):

- sound: rotation of  $\hat{\mathbf{n}}$  and  $\hat{\mathbf{m}}$  around  $\hat{\mathbf{l}}$
- 2 spin waves: motion of  $\hat{\mathbf{d}}$
- 2 orbital waves: motion of  $\hat{\mathbf{l}}$

Higgs modes (13):

- 6 Clapping modes
- 4 Flapping (spin-orbit) modes
- 1 Pseudo-sound mode
- 2 Pseudo-spin modes

$^3\text{He-B}$



Goldstone modes (4):

- sound: motion of  $\varphi$
- 3 spin wave modes: motion of  $\hat{\mathbf{n}}$  and  $\theta$

Higgs modes (14):

- 4 Pair-breaking modes
- 5 Real squashing modes
- 5 Imaginary squashing modes

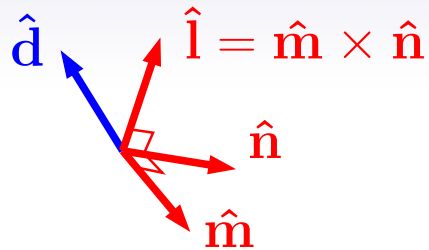
Finate temperature:

- first and second sound

Non-hydrodynamic regime:

- zero sound
- collisionless spin waves

$^3\text{He-A}$



Goldstone modes (5):

- sound: rotation of  $\hat{n}$  and  $\hat{m}$  around  $\hat{l}$
- 2 spin waves: motion of  $\hat{d}$
- 2 orbital waves: motion of  $\hat{l}$

Higgs modes (13):

- 6 Clapping modes
- 4 Flapping (spin-orbit) modes
- 1 Pseudo-sound mode
- 2 Pseudo-spin modes

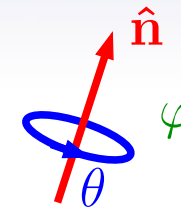
Finate temperature:

- first and second sound

Non-hydrodynamic regime:

- zero sound
- collisionless spin waves

$^3\text{He-B}$



Goldstone modes (4):

- sound: motion of  $\varphi$
- 3 spin wave modes: motion of  $\hat{n}$  and  $\theta$

Higgs modes (14):

- 4 Pair-breaking modes
- 5 Real squashing modes
- 5 Imaginary squashing modes

Rotation of the order parameter changes total spin of the system.

We can write Hamiltonian equations using spin  $S$  and  $R_{\mu j}$  as coordinates.

Energy of spin in magnetic field:

$$F_M = -(\gamma \mathbf{S} \cdot \mathbf{H}) + \frac{\gamma^2 \mathbf{S}^2}{2\chi} \quad (\text{minimum at } \gamma \mathbf{S} = \chi \mathbf{H})$$

Energy of weak spin-orbit interaction in Cooper pairs:

$$F_{SO} = \frac{\chi \Omega_B^2}{15\gamma^2} \sum_{k,j} (R_{jj} R_{kk} + R_{jk} R_{kj}) = \frac{\chi \Omega_B^2}{15\gamma^2} \frac{1}{2} (4 \cos \theta + 1)^2 \quad (\text{minimum at } \theta \approx 104^\circ)$$

Gradient energy:

$$F_{\nabla} = \langle \text{some long expression} \rangle$$

# Leggett equations, texture and spin waves

Leggett equations:

$$\dot{S}_a = [\mathbf{S} \times \gamma \mathbf{H}]_a + T_a(R),$$

$$\dot{R}_{aj} = e_{abc} R_{cj} \left( \frac{\gamma^2}{\chi_B} \mathbf{S} - \gamma \mathbf{H} \right)_b,$$

Gradient energy

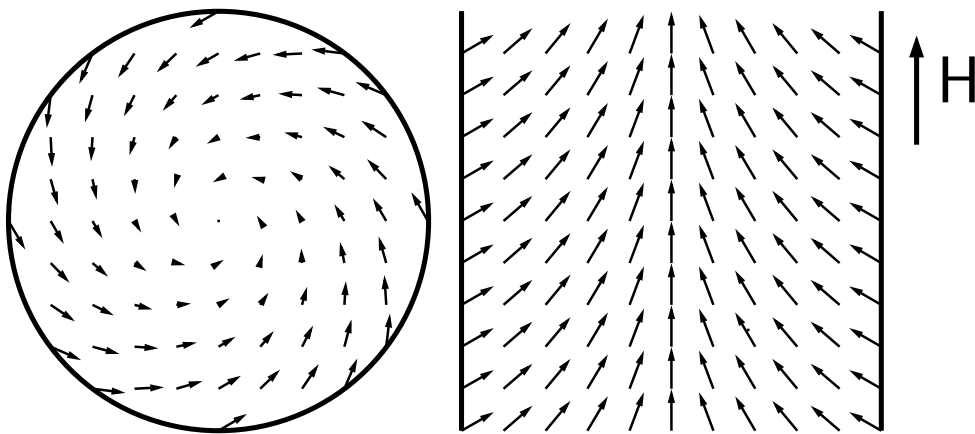
Spin-orbit interaction,  $\Omega_B$

Interaction with walls

Equilibrium distribution of  $R(\hat{\mathbf{n}}, \vartheta)$  – texture.

Motion of  $R(\hat{\mathbf{n}}, \vartheta)$  – spin waves.

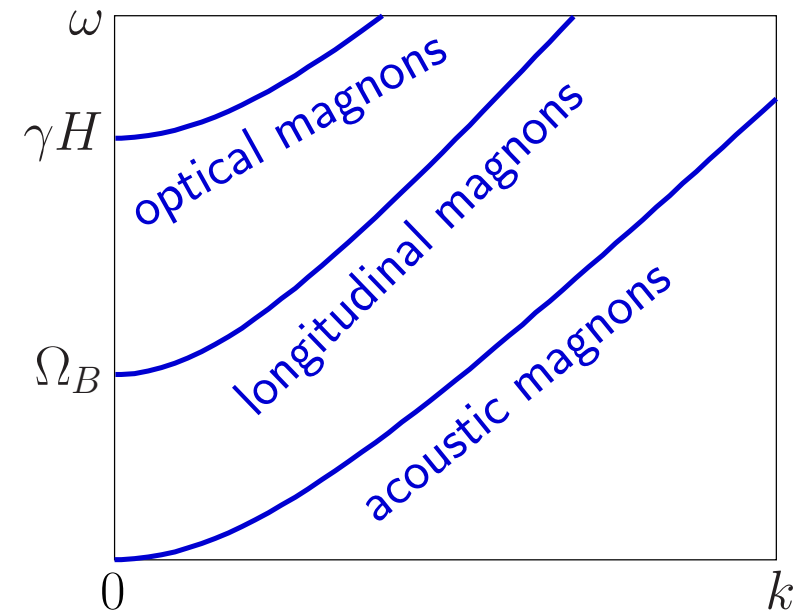
Flare-out texture in a cylindrical cell



$$\vartheta = 104^\circ$$

(Leggett angle, minimum of spin-orbit interaction)

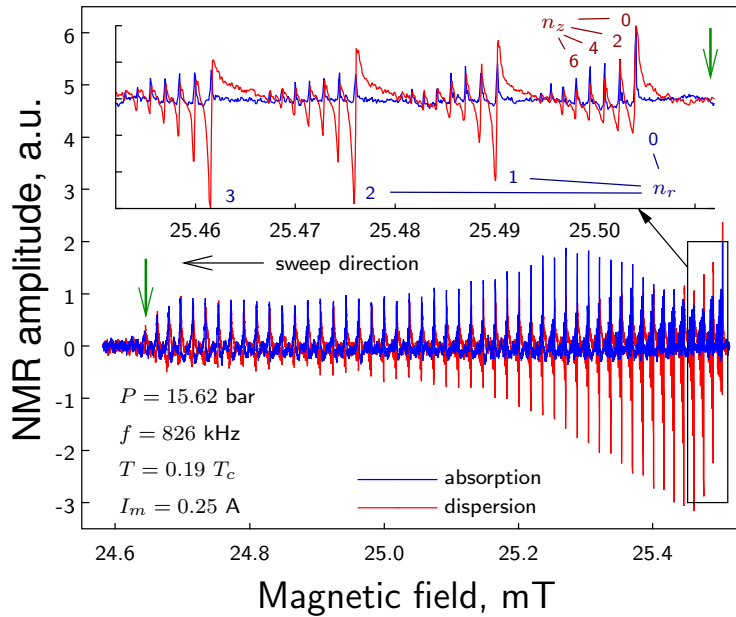
Linear NMR,  $\Omega_B \ll \gamma H$ ,  $\hat{\mathbf{n}} \parallel \mathbf{H}$



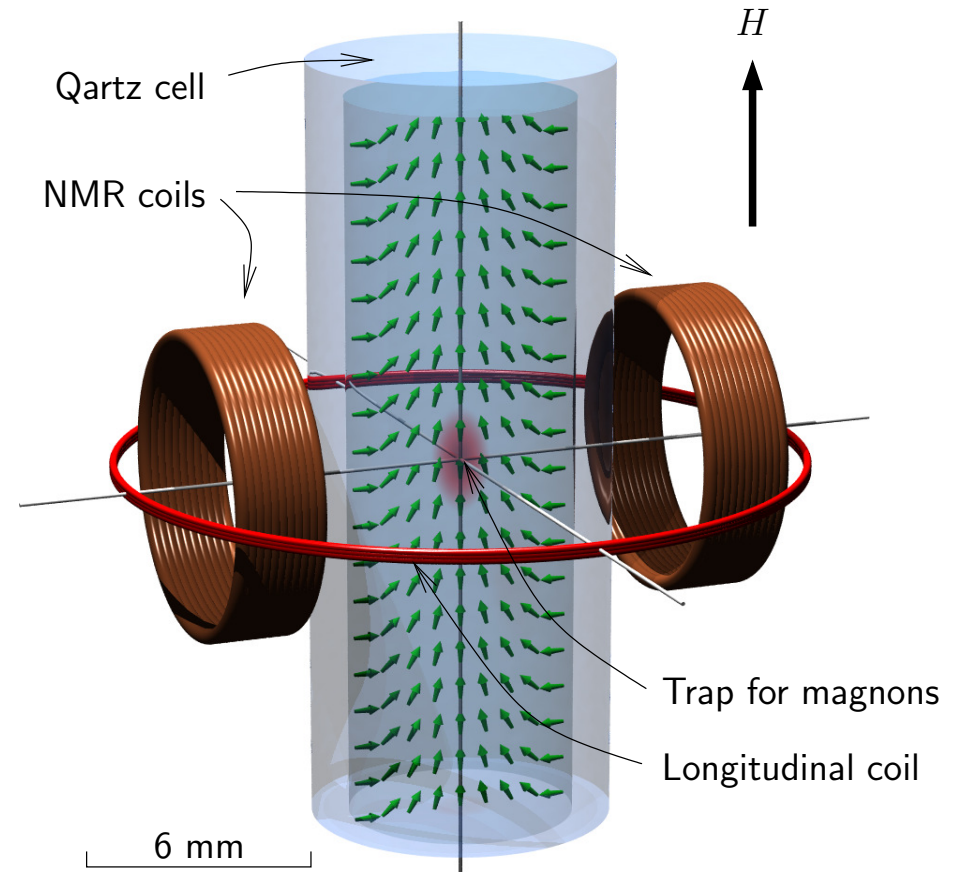
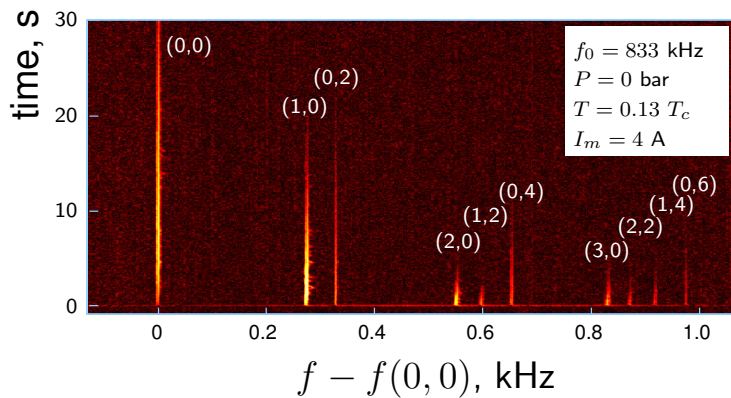
# Trap for magnons

Experiments in Helsinki 2012-2015

## continuous pumping



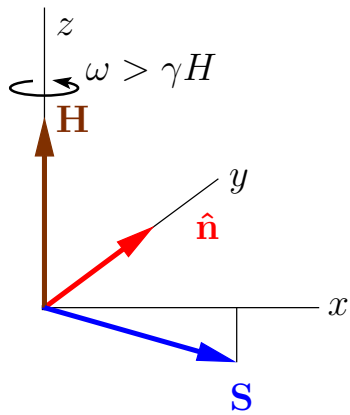
## pulsed pumping



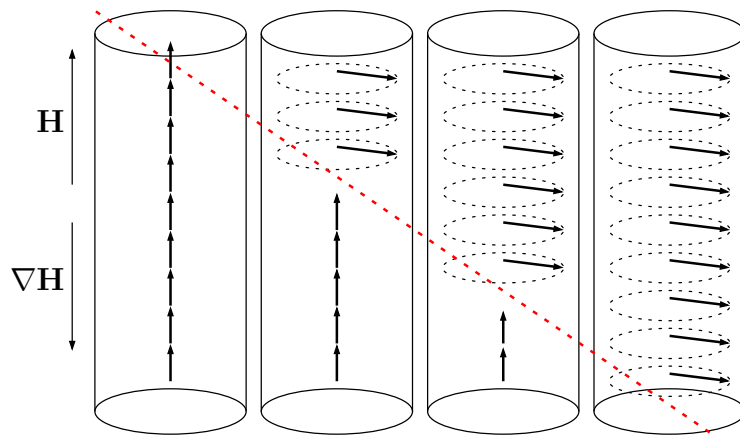
Thermometry, interaction with free surface, with vortices, Studying properties of  $^3\text{He}$ .

# Homogeneously precessing domain (HPD)

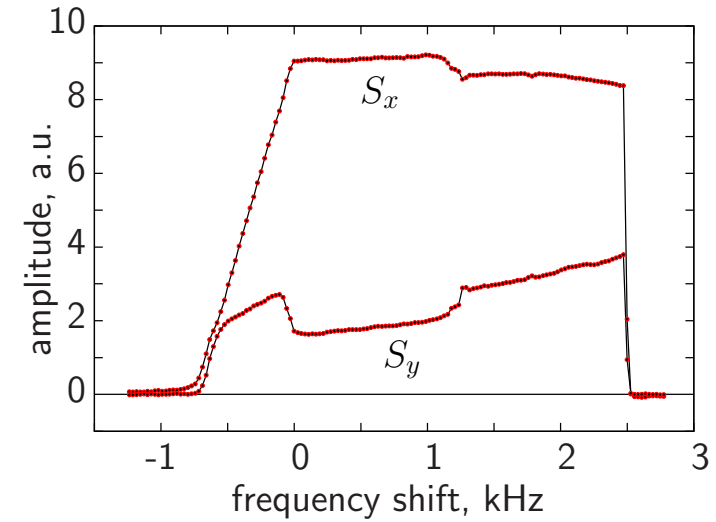
HPD



creation of HPD in CW NMR



$$\omega = \gamma H$$



Experiment with two HPDs (Moscow, 1987)

