



Vibrating wire thermometry in superfluid He³-B

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Royal Holloway, 10.5.2023



Vibrating wire in magnetic field



Wire loop in magnetic field \mathbf{B} is driven by current I and moving with velocity \mathbf{v} .

Force acting on a piece of wire $d\mathbf{l}$: $d\mathbf{F} = I[d\mathbf{l} \times \mathbf{B}]$

EMF voltage:

 $dV = \mathbf{v} \cdot [d\mathbf{l} \times \mathbf{B}]$

By integrating along the wire we have total force and voltage: F = ILB, V = vLB, where v is average velocity

Equation of motion (linear dampling):

$$\ddot{x} = -\omega_0^2 x - \delta \, \dot{x} + \frac{ILB}{m_w} \cos \omega t$$

Using complex notation:

velocity:
$$\mathbf{v} = \frac{i\omega \mathbf{I}LB/m}{\omega_0^2 - \omega^2 + i\omega\delta}$$
 voltage: $\mathbf{V} = \frac{i\omega \mathbf{I}(LB)^2/m}{\omega_0^2 - \omega^2 + i\omega\delta}$









Non-linear damping in He³-B



Classical gas





Amount of colliding particles at both sides per unit time: $\dot{N}_{\pm} = -nA(\langle u \rangle \pm v)$

Momentum transfered by a particle: $p_{\pm} = \pm 2m \langle u \rangle$

Force: $F = p_+ \dot{N}_+ + p_- \dot{N}_- = -4nAm\langle u \rangle v$ Linear damping!

Superfluid ³He-B

Energy spectrum of Bogoliubov quasiparticles:



Normal scttering can be prohibited by conservation of energy: quasiparticle can not loose energy and go below the gap

Force in a simple 1D model:

$$F = -F_0 v_0 (1 - \exp(-v/v_0)) \qquad v_0 = \frac{kT}{p_F} \qquad F_0 = A p_F^2 v_F N(0) \exp\left(-\frac{\Delta}{kT}\right)$$



Non-linear motion of the wire



Non-linear equation of motion:

$$\ddot{x} = -\omega_0^2 x - \bigotimes_0^2 x + \frac{ILB}{m_w} \cos \omega t$$

$$\delta v_0 f(\dot{x}/v_0) \qquad \qquad f(x) \approx x \text{ at } x \to 0$$

1D scattering model: $f(x) = \text{sign}(x)(1 - \exp(-|x|))$

After using van del Pol transformation and averaging over period:

complex velocity: $\mathbf{v} = \frac{i\omega \mathbf{I}LB/m}{\omega_0^2 - \omega^2 + i\omega S(|v|/v_0)\delta}$ voltage: $\mathbf{V} = \frac{i\omega \mathbf{I}(LB)^2/m}{\omega_0^2 - \omega^2 + i\omega S(|V|/V_0)\delta}$ $S(x) = -\frac{2}{x} \int_{0}^{2\pi} f(-x\sin(\omega t) \sin \omega t \frac{d(\omega t)}{2\pi})$ 1.0 0.8 If $f(x) = x + a \operatorname{sign}(x) x^2$ then $S(x) = 1 + \frac{8a}{3\pi} x$ S(x)0.6 0.4 for 1D scattering model: $S(x) = \frac{2}{x} \left(I_1(x) - L_{-1}(x) + \frac{2}{\pi} \right)$, 0.2 0.0 1 where $I_n x$ is modified Bessel function of first kind and $L_n(x)$ is modified Struve function.





Measuring function S



Tracking mode:

1. measure resonance by sweeping frequency at constant drive current I_0

$$\mathbf{V} = \frac{\mathbf{a} I_0 i\omega}{\omega_0^2 - \omega^2 + i\omega S(V)\delta} + \mathbf{b} I_0$$

- 2. Find ${\bf a}$ and ${\bf b}$
- 3. Measure voltage as a function of drive at constant frequency, solve for S.

$$S(V)\delta = \Re\left(\frac{\mathbf{a}\,I}{V-\mathbf{b}\,I}\right)$$

4. Iterate from step 1 with obtained function S





Building model for function S





To process thermometry data we need function S(v). It depends on pressure and temperature (via v_0), on wire intrinsic damping (which depends on magnetic field).

Simultaneous fit of a few frequency sweeps using function S(x) obtained from the left picture



Processing non-linear data





Fitting frequncy sweeps with and without non-linear correction

Applying non-linear correction to tracking mode





Knowledge of function S can give us "zero-velocity" damping δ . It should be converted to temperature.

Theory:

$$\delta = \frac{p_F^2 v_F N(0)}{\rho_w d_w} \exp\left(-\frac{\Delta}{kT}\right)$$

Pre-factor in this formula can be affected by multiple things:

- type of scattering and surface of the wire
- finite coherence length in ³He (for small wire sizes)
- dust on the wire

Possible approach: use 0.127mm tantalum wires for calibrating thin wires:

easier to make and measure wire geometry, less affected by dust, and coherence length effects.